In the beginning, God said



Wave Equations (source free, uniform medium) $\nabla^2 \overline{E} = \mu \epsilon \frac{\partial^2 \overline{E}}{\partial t^2} \quad \nabla^2 \overline{H} = \mu \epsilon \frac{\partial^2 \overline{H}}{\partial t^2}$ *Light is EM waves!*

Then, there was light!

 $\begin{array}{ll} \text{Material} \\ \text{Parameters} \end{array} \begin{cases} \epsilon : \text{permittivity} \\ \mu : \text{permeability} \end{cases}$

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Solutions for Wave Equations

$$\nabla^{2}\overline{\mathbf{E}} = \mu \varepsilon \frac{\partial^{2}\overline{\mathbf{E}}}{\partial t^{2}} \qquad \nabla^{2}\overline{\mathbf{H}} = \mu \varepsilon \frac{\partial^{2}\overline{\mathbf{H}}}{\partial t^{2}}$$

(Plane-wave solutions)

$$\overline{\mathbf{E}} = \overline{\mathbf{x}} \mathbf{E}_0 \mathbf{e}^{\mathbf{j}(\omega \mathbf{t} - \mathbf{k}z)} \quad \text{and} \quad \overline{\mathbf{H}} = \overline{\mathbf{y}} \mathbf{H}_0 \mathbf{e}^{\mathbf{j}(\omega \mathbf{t} - \mathbf{k}z)}$$
$$\omega = 2\pi \mathbf{f} \qquad k = \frac{2\pi}{\lambda} = \omega \sqrt{\mu\varepsilon} \qquad \frac{\mathbf{E}_0}{\mathbf{H}_0} = \sqrt{\frac{\mu}{\varepsilon}} = \eta \qquad [\Omega]$$

Direction of E, H fields? Direction of propagation? Speed of propagation?

$$f \cdot \lambda = \frac{\omega}{k} = \frac{1}{\sqrt{\mu\varepsilon}}$$

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How does the plane-wave solution look like?

For physical representation, $\operatorname{Re}\left[\overline{x}E_{0}e^{j(\omega t - kz)}\right] = \overline{x}E_{0}\cos(\omega t - kz)$



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How does the plane-wave solution look like?



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Propagation Directions

+z direction:
$$e^{j(\omega t - kz)}$$
 or $e^{-j(\omega t - kz)}$
-z direction: $e^{j(\omega t + kz)}$ or $e^{-j(\omega t + kz)}$
+y direction: $e^{j(\omega t - ky)}$ or $e^{-j(\omega t - ky)}$

Arbitarary direction,
$$\overline{k} = \overline{x}k_x + \overline{y}k_y + \overline{z}k_z$$
 with $\overline{R} = \overline{x}x + \overline{y}y + \overline{z}z$
 $e^{j(\omega t - \overline{k} \cdot \overline{R})} = e^{j\omega t}e^{-jk_x x}e^{-jk_y y}e^{-jk_z z}$

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Polarization: Change of E-field direction with time

Linear Polarization $\overline{E} = (\overline{x}E_0 + \overline{y}E_0)e^{j\omega t}e^{jkz}$



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Circular Polarization

$$\overline{E} = \left(\overline{x}E_0 + \overline{y}jE_0\right)e^{j\omega t}e^{jkz}$$



Handedness?

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Elliptical Polarization

$$\overline{E} = \left(\overline{x}E_0 + \overline{y}j2E_0\right)e^{j\omega t}e^{jkz}$$



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Characteristics of media:

$$\begin{cases} \varepsilon : \text{permittivity} \\ \mu : \text{permeability} \end{cases} \begin{cases} \overline{D} = \varepsilon \overline{E} \\ \overline{B} = \mu \overline{H} \end{cases}$$

- Assume $\mu = \mu_0$ in this course.
- With different ε (dielectric media), how do plain-wave solutions chnage?

For example, $\overline{E} = \overline{x}E_0e^{j(\omega t - kz)}$

$$k = \omega \sqrt{\mu \varepsilon} = nk_0 \ (k_0 = \omega \sqrt{\mu \varepsilon_0}, \ n = \sqrt{\frac{\varepsilon}{\varepsilon_0}}; \ \text{refractive index})$$

=> change in λ , phase velocisty $(v_p = \frac{\omega}{k})$, group velocity $(v_g = \frac{\partial \omega}{\partial k})$

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$$\overline{E} = \overline{x}E_0 e^{-j\beta z} e^{-\alpha z} = \overline{x}E_0 e^{-j(\beta - j\alpha)z}; \text{ complex } k$$
(With σ , $\varepsilon_c \equiv \varepsilon - j\frac{\sigma}{\omega}$ and $k = \omega\sqrt{\mu\varepsilon_c}$)

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- Birefringent Media: Change in polarization

Consider $\vec{E}_{in} = (\vec{x} + \vec{y})E_0 e^{j(\omega t - k_0 z)}$ Incident on a birefringent material (different refractive indices for different directions)



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